# Electron energy and electric field estimates in sprites derived from ionized and neutral $N_2$ emissions

J. Morrill,<sup>1</sup> E. Bucsela,<sup>2</sup> C. Siefring,<sup>3</sup> M. Heavner,<sup>4</sup> S. Berg,<sup>5</sup> D. Moudry,<sup>6</sup> S. Slinker,<sup>3</sup> R. Fernsler,<sup>3</sup> E. Wescott,<sup>6</sup> D. Sentman,<sup>6</sup> and D. Osborne<sup>6</sup>

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[1] During the EXL98 aircraft mission, sprites and blue jets were observed by narrow band cameras that measure the  $N_2^+$  1NG (0,1) band at 4278Å and the  $N_2$  2PG (0,0) band at 3370Å. We discuss the observations ( $\sim$ 1 km resolution), instrumental and atmospheric corrections, and altitude profiles of ionized (1NG) and neutral (2PG) emission observed during a specific sprite. The ratio of ionized-to-neutral emission indicates a relative enhancement of ion emission below 55 km. Characteristic electron energies  $(E_{Ch})$  and electric fields (E) are derived from these emission ratios using excitation rates computed from a model that solves the Boltzmann equation as a function of electric field. Up to 55km E follows the breakdown field  $(E_k)$  and  $E_{Ch}$  is ~2.2eV. Above 55 km E drops below  $E_k$  and  $E_{Ch}$  drops to ~1.75eV near 60km. INDEX TERMS: 0342 Atmospheric Composition and Structure: Middle atmosphere-energy deposition; 3304 Meteorology and Atmospheric Dynamics: Atmospheric electricity; 3324 Meteorology and Atmospheric Dynamics: Lightning; 3334 Meteorology and Atmospheric Dynamics: Middle atmosphere dynamics (0341, 0342)

## 1. Introduction

[2] The EXL98 Aircraft Campaign was a NASA sponsored mission to observe sprites and related phenomena from aircraft during July of 1998. A major goal was to measure neutral and ionized molecular nitrogen emission to examine ionization in sprites. The primary instrument configuration included narrow-band imagers that made video observations of the  $N_2^+$  1NG (0,1) band at 4278 and  $N_2$  2PG (0,0) band at 3370Å. These imagers were provided by the Naval Research Laboratory and the University of Alaska Fairbanks, respectively. Other EXL98 observations measured UV sprite spectra [*Heavner et al.*, 2000]. Both observations are affected by atmospheric attenuation [*Morrill et al.*, 1998], and the EXL98 mission minimized this effect by using an aircraft as a high altitude observing platform.

[3] Spectral analysis studies indicate the presence of  $N_2^+$  Meinel emissions in sprites [*Green et al.*, 1996; *Morrill et al.*, 1998]. These identifications involve relatively weak signals, as expected given  $N_2^+$  (*A*) quenching [*Piper et al.*, 1985]. *Green et al.* [1996] compare model results with sprite spectra and estimate electron energies near 1 eV. Model estimates of blue and near-UV sprite emissions of  $N_2^+$  and  $N_2$  as well as electron energetics in sprites are presented by

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numerous groups [Fernsler and Rowland, 1996; Pasko et al., 1997; Rowland, 1998].

[4] Photometric observations of  $N_2^+$  and  $N_2$  in the blue and UV have been made from the ground [*Armstrong et al.*, 1998; *Suszcynsky et al.*, 1998; *Armstrong et al.*, 2000] hereafter (I, II, III). These studies produce precise time histories of sprite emissions with photometers. However, they integrate over a range of altitudes so variations in the time history with altitude were not determined. Initial work (I, II) used wide filter passbands that allowed both ionized (1NG) and neutral (2PG) emission into each channel and required extensive spectral modeling. Later work (III) used smaller passbands but still integrated over 15–25 km. The present observations integrate the sprite emission over 33 ms so that temporal information is limited. However, aircraft-based imaging and narrow-band filters provide altitude resolved profiles of the 1NG and 2PG emission with resolution improved over previous studies.

[5] In this paper we discuss the narrow-band EXL98 observations and the instrumental and atmospheric corrections used in the analysis. The resulting ratios of  $N_2^+$  1NG (0,1) to  $N_2$  2PG (0,0) emission yield excitation rate ratios of  $N_2^+$  (*B*) to  $N_2$  (*C*). Characteristic electron energy ( $E_{Ch}$ ) and electric field (*E*) estimates come from a comparison of the observed excitation rates with rates determined from a model that solves the Boltzmann equation as a function of electric field.

# 2. Observations

[6] Two sprites were observed at ~0900 UT on July 19, 1998, over Wisconsin. These sprites were observed from a single location so altitude and range information is based on the location of the two associated lightning strokes and known sprite morphology. Figure 1a shows the image from a broadband camera. Range to the sprite is ~ 360 km and the aircraft altitude was 14 km.

[7] Figure 1a also shows the field-of-view (FOV) of the 1NG and 2PG filtered cameras (inserted box). Figures 1b and 1c show the emission of the 1NG (0,1) and 2PG (0,0), respectively. The broadband image is saturated and not useful for this analysis. The central wavelengths and FWHM band-widths for these filters are 4282(14)Å and 3410(72)Å, respectively. The 2PG (0,0) is in the blue wing of the 3410Å filter. The two possible contaminating bands are the 2PG (1,5) and (4,5) bands in the 4282 and 3410Å passbands, respectively. Given filter rejection and intensity estimates, neither band produces a significant signal. 1NG observations used a Dage/MTI SIT VE1000 camera with a Fuji f/0.7 lens. 2PG observations used a GEN IIUV V53-1845 Video Scope camera and a Lyman Alpha II f/1.7 lens. This analysis focuses on the left sprite in Figure 1a.

## 3. Analysis and Results

[8] These observations require a number of corrections. The 2PG image is scaled to the 1NG image. Figures 1b and 1c show three narrow vertical regions with the left sprite in the central region. The two outer regions provide the background. The central region is summed horizontally and the residual sprite emission is the difference between the smoothed fits to the sprite and back-

<sup>&</sup>lt;sup>1</sup>E. O Hulburt Center for Space Research, Naval Research Laboratory, Washington, DC, USA.

<sup>&</sup>lt;sup>2</sup>Goddard Earth Science and Technology Center, Baltimore, MD, USA.
<sup>3</sup>Plasma Physics Division, Naval Research Laboratory, Washington DC, USA.

<sup>&</sup>lt;sup>4</sup>Los Alamos National Laboratory, Los Alamos, NM, USA.

<sup>&</sup>lt;sup>5</sup>Computational Physics Inc., Springfield, VA, USA.

<sup>&</sup>lt;sup>6</sup>Geophysical Institute, University of Alaska Fairbanks, Fairbanks, AL, USA.



Figure 1. Image of sprites from (a) broad-band camera, (b) 1NG filtered camera (c) 2PG filtered camera.

ground profiles. The 1NG and 2PG profiles are in Figure 2 and have  ${\sim}1$  km altitude resolution.

[9] Other instrumental and atmospheric corrections are applied to these profiles. Corrections account for the peak wavelength (angle dependent) and spectral shape of the passband filters, camera flat field, absolute calibration and altitude-dependent atmospheric attenuation (MOSART atmospheric transmission model [*Morrill et al.*, 1998]). All corrections are slowly varying with altitude so large changes in the intensity profiles cannot be attributed to instrumental or atmospheric effects.

[10] Figure 3 shows the corrected intensity profiles for the ionized (1NG) and neutral (2PG) emissions. The 1NG emission remains roughly constant at all altitudes with a slight increase at lower altitudes. The 2PG profile is brighter and more structured than the 1NG profile. Above 55 km the neutral emission dominates



**Figure 2.** Residual sprite emission profiles for the 1NG (0, 1) and 2PG (0, 0) band emissions.

and varies significantly while below 55 km it is roughly constant. The ratio of these emissions appears in Figure 4 and shows two regions; the upper portion or "body" and the lower portion or "tendrils." These results indicate a relative enhancement in ionization in the tendrils.

[11] A number of assumptions allow  $E_{Ch}$  and E to be estimated from the intensity ratios. The intensity is

$$I_{\nu'\nu''} = k_e N_2 n_e A_{\nu'\nu''} q / \left( 1 / \tau_{\nu'} + k_q^{N_2} N_2 + k_q^{O_2} O_2 \right)$$
(1)

with  $k_e$  the excitation coefficient,  $N_2$  and  $n_e$  the nitrogen and electron densities,  $A_{\nu\nu\nu'}$  the transition probability, q the Franck-Condon factor,  $\tau_{\nu'}$  the lifetime, and  $k_q$  the quenching coefficients (see Table 1).

[12] Equation (1) applies to the instantaneous intensities, but our emission data are time-integrated. To allow for this difference we assume: a sprite consists of waves of ionization (i.e., streamers);



Figure 3. 1NG (0, 1) and 2PG (0, 0) intensity profiles.



Figure 4. 1NG (0, 1)/2PG (0, 0) intensity ratios.

 $E_{Ch}$  is determined by the instantaneous electric field E;  $E_{Ch}$  is small compared with the threshold excitation energies of the 1NG and 2PG upper states (~10–20 eV). In this case,  $k_e$  is a strong function of E. Ahead of the streamer E and  $k_e$  are large while  $n_e$  is small. In the streamer head E and  $k_e$  are large with  $n_e$  increasing. Finally, in the streamer body E and  $k_e$  are small while  $n_e$  is large [*Raizer et al.*, 1998]. Thus,  $I_{\nu'\nu'}$  peaks in the streamer head where both  $k_e$  and  $n_e$ are large. For the case where the emission peaks near the streamer head and persists for less time than for the streamer to cross one pixel ( $\Delta h \sim 1 \text{ km}$ ) [*Stanley et al.*, 1999], the ratio of the peak intensities is approximately the ratio of the time-integrated emissions. We then estimate  $E_{Ch}$  and E in the streamer head as a function of altitude using equation (1), the measured intensity ratio, and the results of a Boltzmann code [*Slinker et al.*, 1990] run with 20%  $O_2/80\% N_2$ .

[13] For the model to apply,  $\Delta h$  should be small compared with the neutral scale height (~10 km) but large compared with the ionization front (~0.1 km [*Gerken et al.*, 2000]). The latter assumption is best met at low altitude, since the width of the front varies inversely with density. If the assumption fails, the model underestimates the values of  $E_{Ch}$  and E in the streamer head, because the lower excitation energy of the 2PG upper state allows its emission to persist (III). The results from this model appear in Figure 5. Errors in  $E_{Ch}$  are ~0.1eV and reflect intensity uncertainties.

#### 4. Discussion

[14] Sprite processes require observations at small spatial, spectral, and temporal scales. So far, no single study has simultaneously probed all three scales. Photometric studies of  $N_2^+$  and  $N_2$  emissions (I, II, III) yield detailed time history but little spatial information. High speed video [*Stanley et al.*, 1999] yields spatial and temporal results but little spectral data. The current results give

Table 1. Radiative and Kinetic Parameters

Parameter	$N_2(C)$	$N_2^+(B)$
$\tau(s)$	$3.71 \times 10^{-8}$	$6.23 \times 10^{-8}$
A $(s^{-1})$	$1.31 \times 10^{7}$	$3.71 \times 10^{6}$
q	$5.45 \times 10^{-1}$	$8.83 \times 10^{-1}$
$\hat{k}_a^{O_2}(cc/s)$	$5.48  imes 10^{-10}$	$6.20 \times 10^{-10}$
$k_q^{q_{N_2}}(cc/s)$	$3.59\times 10^{-11}$	$7.50  imes 10^{-11}$



**Figure 5.** Electric field (E) and electron energy  $(E_{Ch})$  from intensities as well as breakdown field  $(E_k)$ .

limited temporal information but yield spatial (altitude) and spectral information. Given the differing 1NG and 2PG emission profiles, neither photometry nor high-speed video can resolve issues of sprite energetics without considering the present results.

[15] The present analysis is valid when the emission ratio is constant in time or dominated by one value. This should be valid for a propagating streamer breakdown, since 1NG/2PG emission intensity is weighted to times where  $E_{Ch}$  and  $n_e$  are large. The 1NG/2PG intensity ratio (Figure 4) shows two distinct regions, above and below 55 km.  $E_{Ch}$  in Figure 5 decreases with decreasing altitude to ~60 km. A transition occurs between 60–55 km with the  $E_{Ch}$  increasing to ~2.2eV below 55 km. Maxwell-Boltzmann (MB) distributions do not accurately yield  $E_{Ch}$  in Figure 5.

[16] We interpret our results in light of high speed video observations [D Moudry, private comm.] that show "tendril" emissions occur for  $\sim 1$  ms while "body" emissions continues for  $\sim 10$  ms. This is consistent with ionization occuring during an initial hot phase ( $\sim 1$ ms) (I, II, III) at 40–70 km with a longer cool phase above 55 km where neutral emission dominates. Below 55 km, one interpretation is that emissions are from the ionizing portion of the streamer tip and so indicate the electric field is capped at the breakdown field. As a result,  $E_{Ch}$  values below 55 km are the characteristic electron energies during ionization. Above 55 km, the measured  $E_{Ch}$  are upper bounds for  $E_{Ch}$  during the cooler phase.

## 5. Conclusion

[17] Analysis of narrow passband sprite observations of 1NG and 2PG emissions yield estimates of  $E_{Ch}$  and E as a function of altitude, using a Boltzmann model. The results show the electric field follows the local breakdown field,  $E_{k}$ , at altitudes below 55 km with  $E_{Ch} \sim 2.2$  eV. Above 55 km, 2PG emission is  $\sim 5 X$  above that at lower altitudes, the derived field is  $\sim 30\%$  below  $E_{k}$ , and  $E_{Ch}$  reaches a minimum of  $\sim 1.75$  eV at 60 km. The latter field values are below  $E_k$  since the model underestimates the peak E and  $E_{Ch}$  during ionization at low gas density. Above 55 km, measured  $E_{Ch}$  represents the upper bound to  $E_{Ch}$  during the cooler phase after ionization. These results indicate that most emission in sprites

results from classical gas breakdown followed, at higher altitudes, by an afterglow with lower  $E_{Ch}$ . This method shows promise for providing time varying  $E_{Ch}$  in sprites when coupled with a dynamic streamer model and spatial, spectral, and temporally resolved data.

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- J. Morrill, C. Siefring, S. Slinker, and R. Fernsler, Naval Research Laboratory, Washington, DC 20375, USA.
- E. Bucsela, NASA/GSFC, Code 916 Greenbelt, MD 20771, USA.
- M. Heavner, NIS-1, Los Alamos National Laboratory Los Alamos, NM 87545, USA.
- S. Berg, Computational Physics, Inc. Springfield, VA 22151, USA.

D. Moudry, D. Osborne, D. Sentman, and E. Wescott, Geophysical Institute, University of Alaska Fairbanks, Fairbanks AK 99775, USA.